

Topology and Homotopy

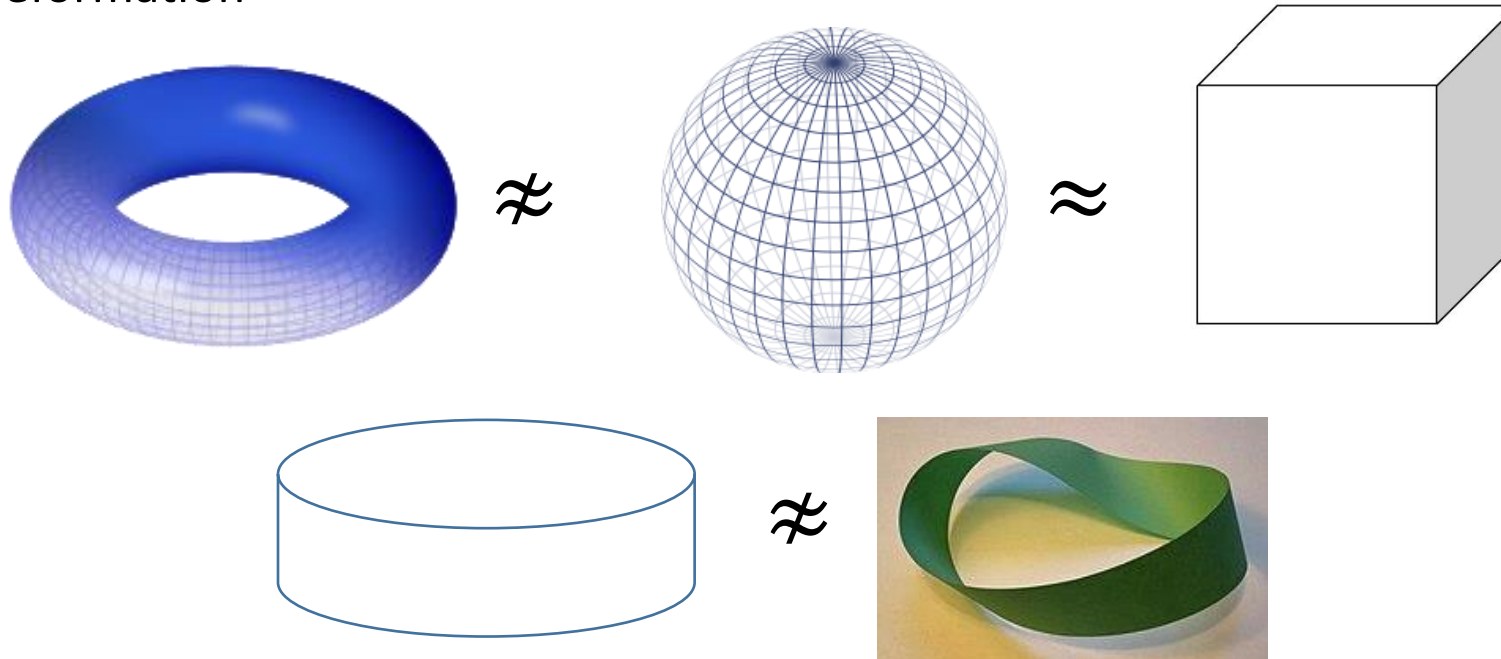
Stanislas Rohart

Laboratoire de Physique des Solides

CNRS/Université Paris-Saclay, France

Introduction: Homeomorphism

Topology aim at classify objects in classes where objects can be transformed one to another using a continuous deformation



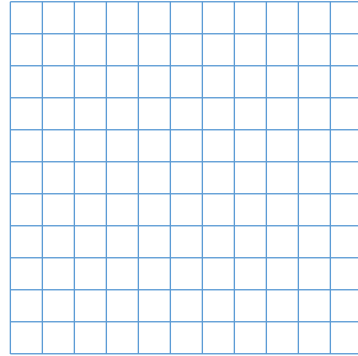
Definition: **Homeomorphism** \approx

Two objects are homeomorph ($S_1 \approx S_2$) if there is a continuous transformation from one to another

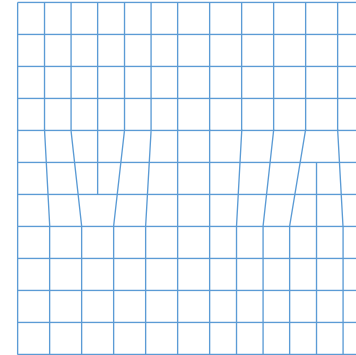
Introduction: Homeomorphism in solid state physics

Example of dislocations

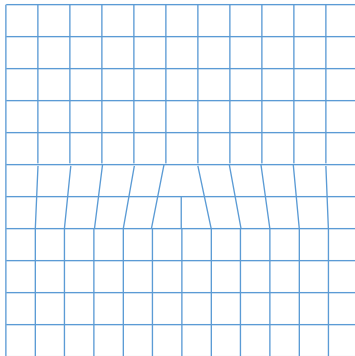
Perfect crystal



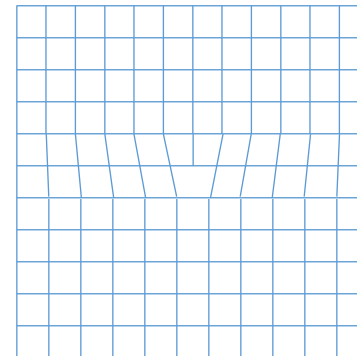
With 2 opposite dislocations
Topology is preserved



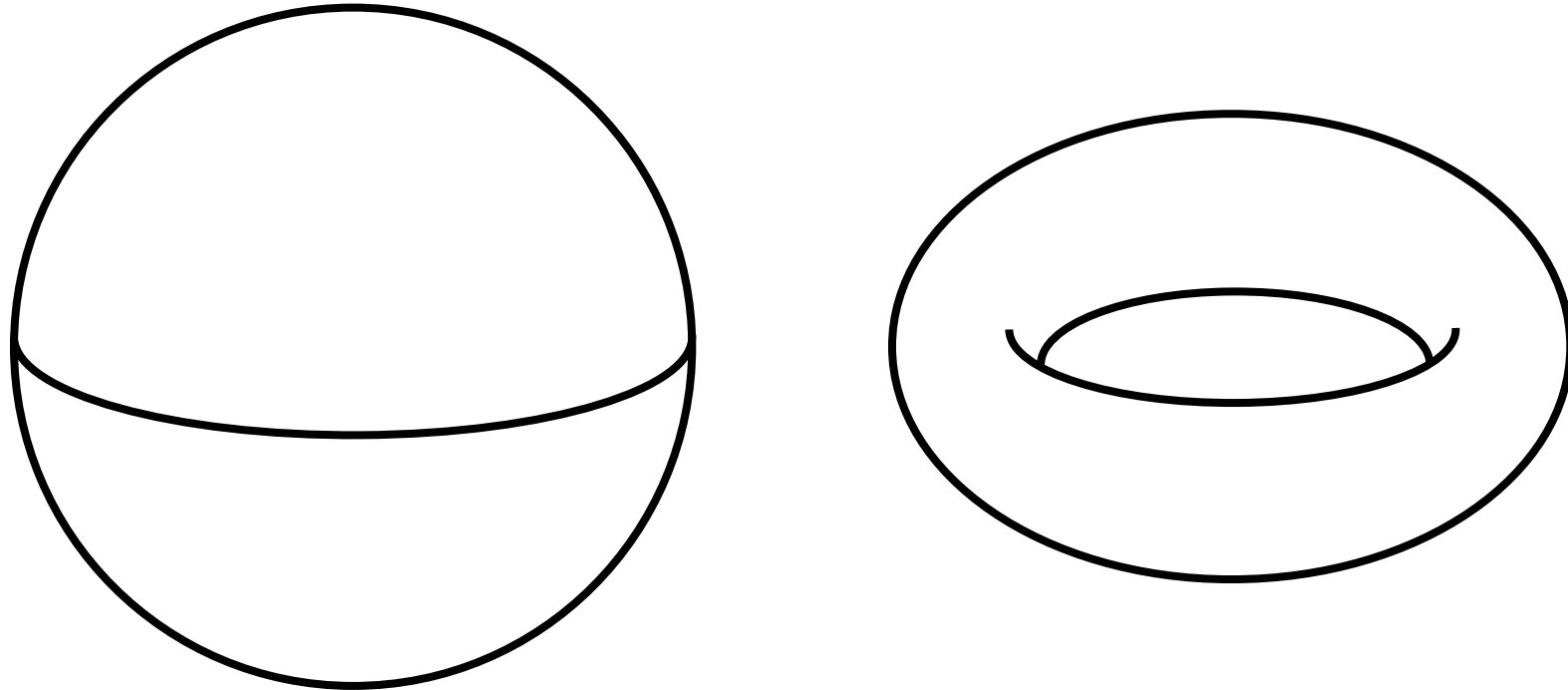
D1 : Crystal with
a dislocation



D2 : Crystal with
another dislocation



Introduction: Objectives of the algebraic topology



How to define formally the topological difference?
How to define formally a hole?
Can we define topologically invariant quantities?

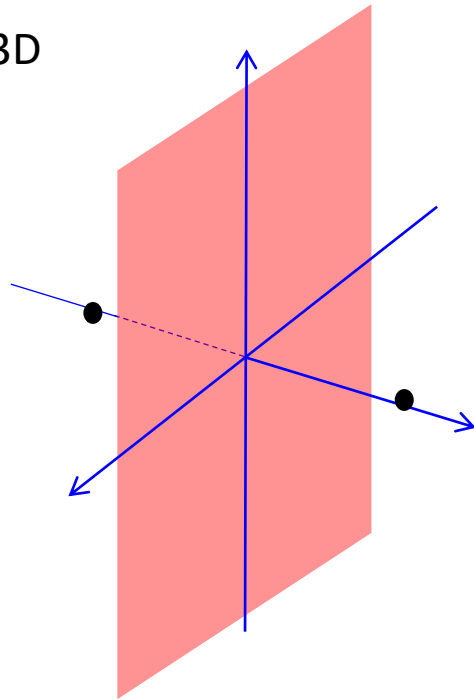
- 1. Introduction: homeomorphism**
- 2. How to catch a topological defect or texture**
- 3. Homotopy and homotopy group**
- 4. Geometrical space and order parameter space**
- 5. Topological defects and topologically stable configurations**

How to catch a topological defect

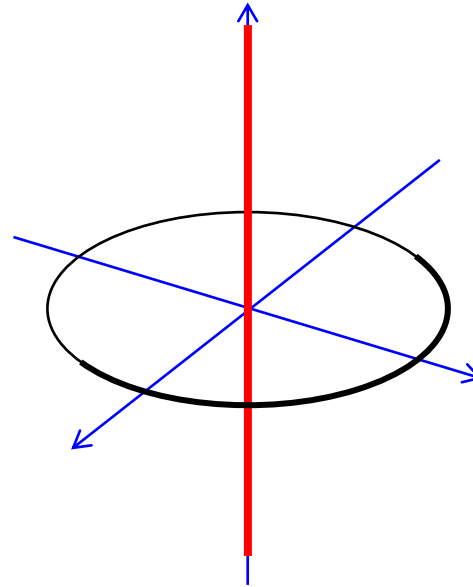
The prey and the hunter rule

A contour surrounds the defect to catch it.

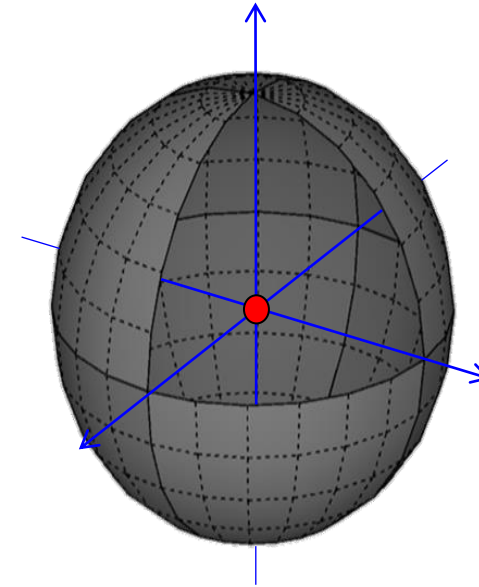
Example in 3D



Surfacic defect



Linear defect



Punctual defect

Hunter's rule

$$d' : \text{Defect dimension} \rightarrow d' + r = d - 1 \leftarrow d : \text{Space dimension}$$

r : Contour dimension

Applications of the hunter's rule to dislocations

$$d' + r = d - 1$$

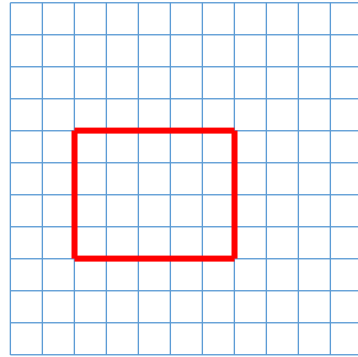
$$d' = 0$$

$$d = 2$$

$$r = 1$$

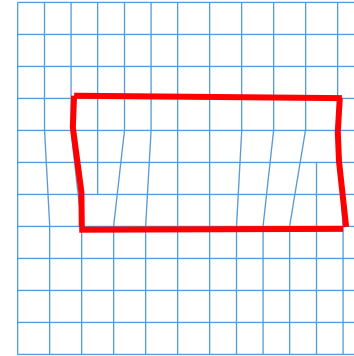
$$\vec{b} = \vec{0}$$

Perfect crystal



≈

With 2 opposite dislocations
Topology is preserved



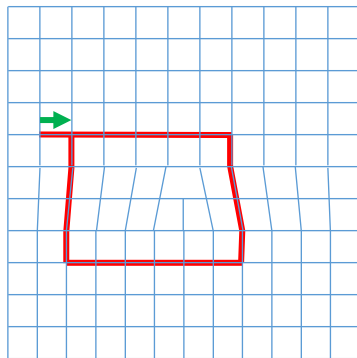
$$\vec{b} = \vec{0} = \vec{b}_{D1} + \vec{b}_{D2}$$

Any contour $(n \times p \times n \times p)$
is closed

D1 : Crystal with
a dislocation

≠

$$\vec{b}_{D1} = a\vec{x}$$

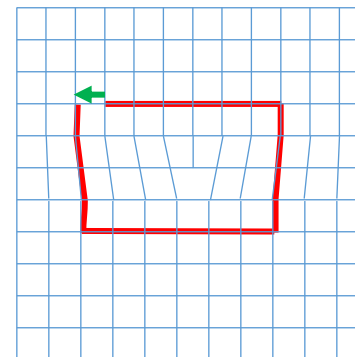


The same contour is not closed

D2 : Crystal with
another dislocation

≠

$$\vec{b}_{D2} = -a\vec{x}$$



The result is different again

How can we
characterize the
dislocation?
-> **Burgers vector**
Topological invariant

- 1. Introduction: homeomorphism**
- 2. How to catch a topological defect or texture**
- 3. Homotopy and homotopy group**
- 4. Geometrical space and order parameter space**
- 5. Topological defects and topologically stable configurations**

Homotopies

1. Properties of contours Explanation for 1D contours γ dans un espace \mathcal{E}

Definition: Path $\gamma_{x_0x_1}$

$\gamma_{x_0x_1}: [0,1] \mapsto \mathcal{E}$ is a continuous application of $[0,1]$ over the space \mathcal{E} so that $\gamma(0) = x_0$ and $\gamma(1) = x_1$

Definition: Closed path or contour γ in x_0

$\gamma: [0,1] \mapsto \mathcal{E}$ is a continuous application of $[0,1]$ over the space \mathcal{E} so that $\gamma(0) = \gamma(1) = x_0$

Definition: Homotopy \sim

γ_1 and γ_2 are homotope ($\gamma_1 \sim \gamma_2$) if there is a continuous transformation from one to another.

Definition: Class of contour $[\gamma]$

$\forall \gamma_1$ and $\gamma_2 \in [\gamma], \gamma_1 \sim \gamma_2$

Definition: Composition .

The composition $\gamma_1 \cdot \gamma_2$ of two contours is a contour. The composition commutes: $\gamma_1 \cdot \gamma_2 \sim \gamma_2 \cdot \gamma_1$

Additional property : $[\gamma_1 \cdot \gamma_2] = [\gamma_1] \cdot [\gamma_2]$

Independence on the starting point

In a continuous space, there is always a path $\gamma_{x_0x_1}$ that links to points. A contour γ starting from x_0 is homotope to a contour $\gamma' = \gamma_{x_1x_0} \cdot \gamma \cdot \gamma_{x_0x_1}$ starting from x_1

Neutral contour A neutral contour c is such that $\forall \gamma, \gamma \cdot c \sim \gamma$.

If $\gamma \sim c$, then γ is neutral.

Reverse contour $\forall \gamma, \exists \gamma^{-1}$ so that $\gamma \cdot \gamma^{-1} = c$.

If $\gamma \sim \gamma^{-1}$ then $\gamma \sim c$

Homotopy group

Definition: Homotopy group

The homotopy group is the group of the classe of paths

Dimension of the contour → $\pi_1(\mathcal{E}) = \{ [\gamma] \}$ ← Ensemble of class of contour in \mathcal{E}

Space or volume

Special case: If $\forall \gamma$ in $\mathcal{E}, \gamma \sim c$ (or $\gamma \in [c]$),
the homotopy group $\pi_1(\mathcal{E}) = \{ [c] \}$ contains a single element.
It is said to be trivial.

Ex: 2D plane $\mathbb{P} = \mathbb{R}^2$: $\pi_1(\mathbb{P})$ is trivial

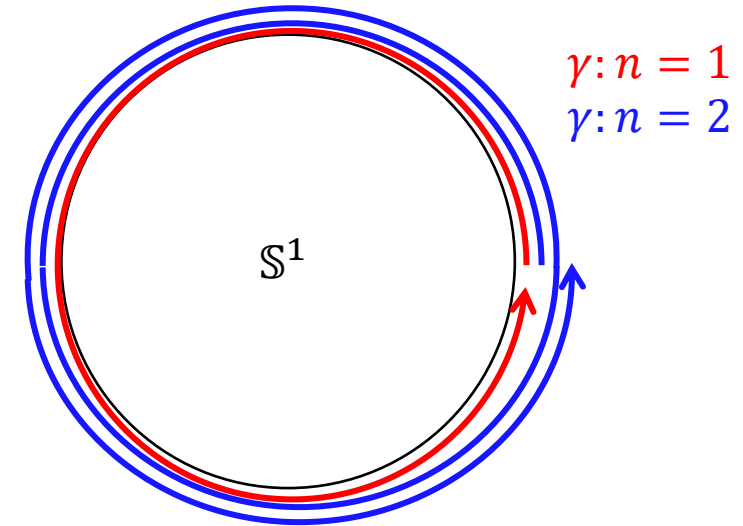
Homotopy group of the circle $\mathcal{E} = \mathbb{C} = \mathbb{S}^1$

Degree $\deg(\gamma)$ of a contour: number of tour around the circle.

$$\deg(\gamma) \in \mathbb{Z}$$

Properties

- $\deg(c) = 0$
- $\deg(\gamma) = -\deg(\gamma^{-1})$
- $\deg(\gamma_1 \cdot \gamma_2) = \deg(\gamma_1) + \deg(\gamma_2)$



The degree is called **topological invariant**
 $n = \deg(\gamma)$

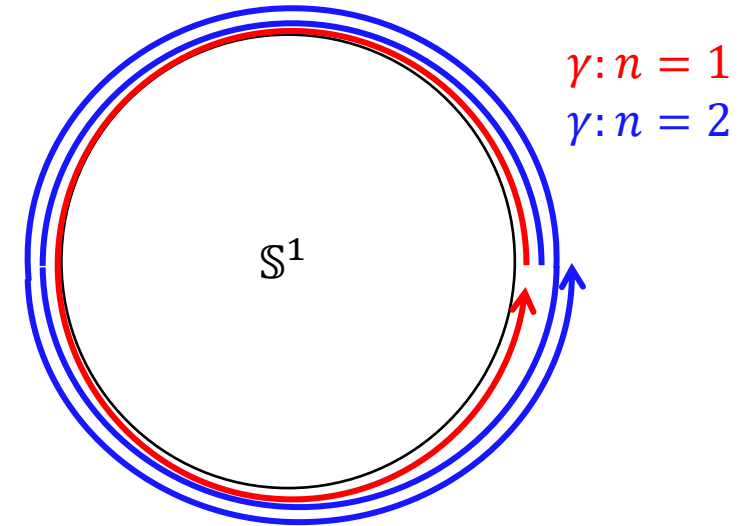
Homotopy group of the circle $\mathcal{E} = \mathbb{C} = \mathbb{S}^1$

Example: $\mathbb{S}^1 = \mathbb{C}_{|z|=1}$

$$z \in \mathbb{S}^1 \rightarrow z = e^{i\theta}$$

Any contour γ is homotope to $\gamma_n(t) = e^{2i\pi nt}$
with $t \in [0,1]$ et $n \in \mathbb{Z}$

$$\text{deg}(\gamma) = n = \frac{1}{2\pi} \oint_{\gamma} dz = \int_0^1 n dt = n$$



Homotopy group of the circle $\mathcal{E} = \mathbb{C} = \mathbb{S}^1$

Any contour is defined by its degree or topological invariant

We can build **the classes of contour** $[\gamma_n]$ such that all the contours of the class share the same degree n .

$$\pi_1(\mathbb{S}^1) = \{ [\gamma_n], n \in \mathbb{Z} \}$$

Since \mathbb{Z} fully defines the homotopy group (isomorphism between $\{ [\gamma_n], n \in \mathbb{Z} \}$ and \mathbb{Z}) the homotopy group is usually written using the group of topological invariants:

$$\pi_1(\mathbb{S}^1) = \mathbb{Z}$$

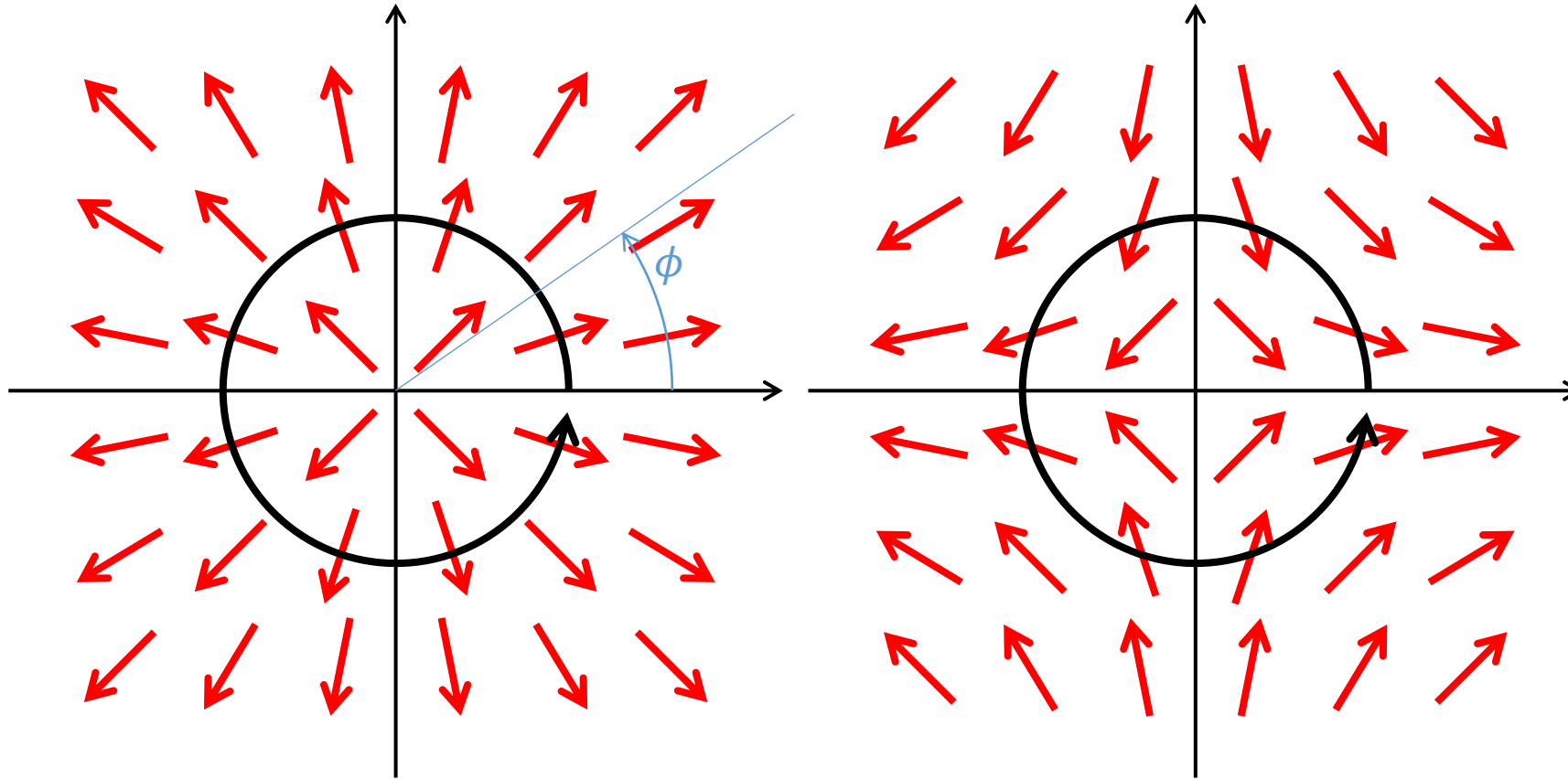
The composition operation between contours becomes an addition in \mathbb{Z}

$$\begin{aligned} \deg(\gamma_1 \cdot \gamma_2) &= \deg(\gamma_1) + \deg(\gamma_2) \\ n(\gamma_1 \cdot \gamma_2) &= n(\gamma_1) + n(\gamma_2) \end{aligned}$$

Homotopy group of the circle $\mathcal{E} = \mathbb{C} = \mathbb{S}^1$

Vortex and antivortex

We consider XY spins that live on \mathbb{S}^1 :
Spins are defined by an angle θ .



Vortex: $\theta = \phi$
 $n = 1$

AntiVortex: $\theta = -\phi$
 $n = -1$

Hunter's rule:

$$d' + r = d - 1$$

The defect is puntual $d' = 0$

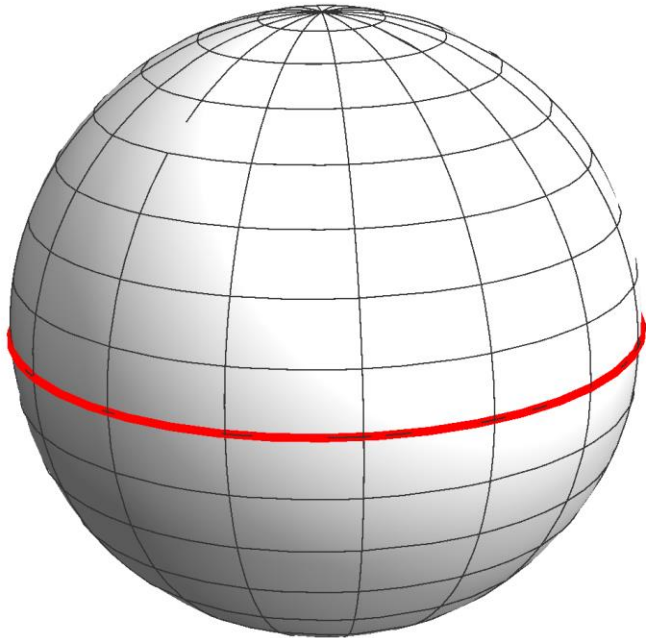
Space is 2D $d = 2$

Contour is $r = 1$ (circle)

Described by an angle ϕ ($0 \rightarrow 2\pi$)

Homotopy group of the sphere S^2

1. First group $\pi_1(S^2)$



Any contour is homotope to a single point (neutral contour):

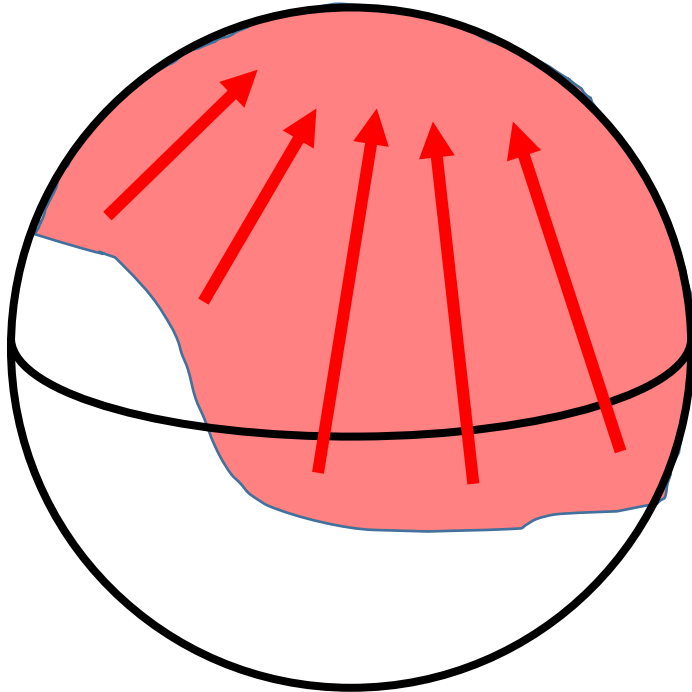
The homotopy group is trivial

$$\pi_1(S^2) = \{c\} = 0$$

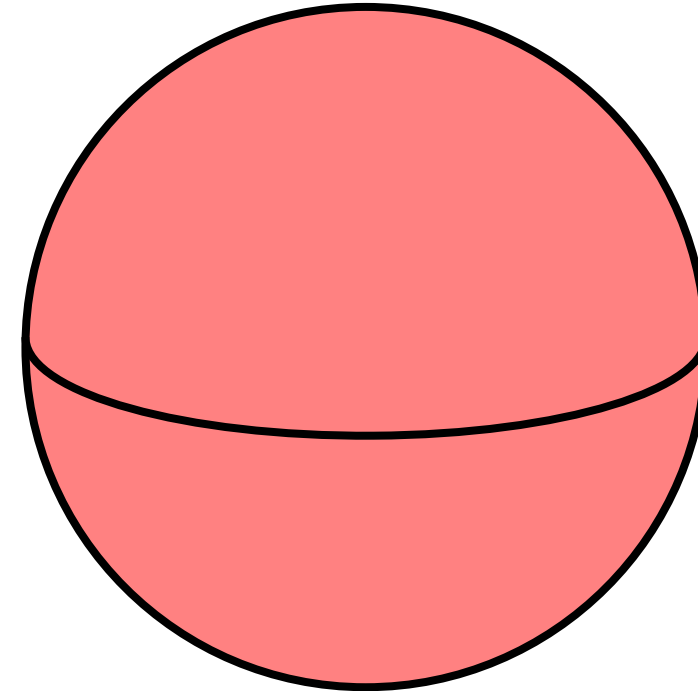
Homotopy group of the sphere S^2

2. Second group $\pi_2(S^2)$

A contour corresponds to a compact surface



This contour is homotope to a point
Trivial or neutral element.



First non trivial element: the sphere is fully covered.

$$\pi_2(S^2) = \mathbb{Z}$$

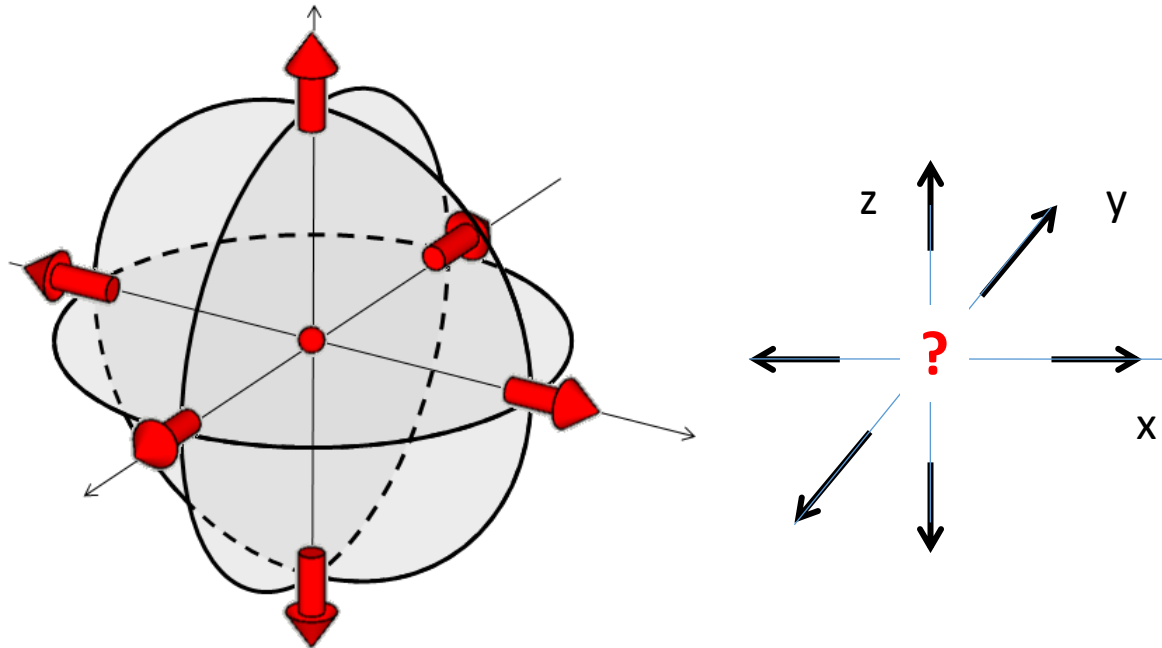
Topological invariant: nb of time the sphere is wrapped.

Homotopy group of the sphere \mathbb{S}^2

2. Second group $\pi_2(\mathbb{S}^2)$

Example of a Bloch point (magnetic monopole)

We consider Heisenberg spin: they live on \mathbb{S}^2 (two angles θ and ϕ)



Hunter's rule:

$$d' + r = d - 1$$

The defect is puntual $d' = 0$

Space is 3D $d = 3$

Contour is $r = 2$ (sphere)

The spin texture found on the spherical contour corresponds to a portion of the \mathbb{S}^2 .

Around the Bloch point, all the spin orientations are found. The spin \mathbb{S}^2 sphere is fully covered.

$$n = \pm 1$$

First description: Feldtkeller Z. Angew. Phys. (1965)

Homotopy group of the spheres S^p

General results:

$$\pi_n(S^p) = 0 \text{ if } n < p$$

$$\pi_n(S^p) = \mathbb{Z} \text{ if } n = p$$

topological invariant: $\frac{1}{|S^p|} \int_{\gamma_n} ds$

(number of time the sphere is covered)

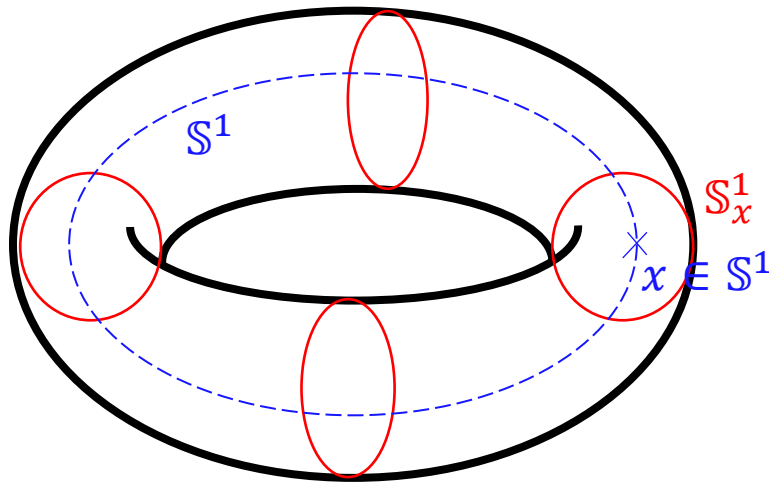
	π_1	π_2	π_3	π_4	π_5	π_6	π_7	π_8	π_9	π_{10}	π_{11}	π_{12}	π_{13}	π_{14}	π_{15}
S^0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
S^1	\mathbb{Z}	0	0	0	0	0	0	0	0	0	0	0	0	0	0
S^2	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{12}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_3	\mathbb{Z}_{15}	\mathbb{Z}_2	\mathbb{Z}_2^2	$\mathbb{Z}_{12} \times \mathbb{Z}_2$	$\mathbb{Z}_{84} \times \mathbb{Z}_2^2$	\mathbb{Z}_2^2
S^3	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{12}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_3	\mathbb{Z}_{15}	\mathbb{Z}_2	\mathbb{Z}_2^2	$\mathbb{Z}_{12} \times \mathbb{Z}_2$	$\mathbb{Z}_{84} \times \mathbb{Z}_2^2$	\mathbb{Z}_2^2
S^4	0	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	$\mathbb{Z} \times \mathbb{Z}_{12}$	\mathbb{Z}_2^2	\mathbb{Z}_2^2	$\mathbb{Z}_{24} \times \mathbb{Z}_3$	\mathbb{Z}_{15}	\mathbb{Z}_2	\mathbb{Z}_2^3	$\mathbb{Z}_{120} \times \mathbb{Z}_{12} \times \mathbb{Z}_2$	$\mathbb{Z}_{84} \times \mathbb{Z}_2^5$
S^5	0	0	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{24}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{30}	\mathbb{Z}_2	\mathbb{Z}_2^3	$\mathbb{Z}_{72} \times \mathbb{Z}_2$
S^6	0	0	0	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{24}	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_{60}	$\mathbb{Z}_{24} \times \mathbb{Z}_2$	\mathbb{Z}_2^3
S^7	0	0	0	0	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{24}	0	0	\mathbb{Z}_2	\mathbb{Z}_{120}	\mathbb{Z}_2^3
S^8	0	0	0	0	0	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{24}	0	0	\mathbb{Z}_2	$\mathbb{Z} \times \mathbb{Z}_{120}$

A. Hatcher, *Algebraic Topology* (Cambridge University Press, Cambridge, 2002 and online available at <http://www.math.cornell.edu/~hatcher>).

Graph source: Wikipedia

Homotopy group of the torus \mathbb{T}^2

The torus corresponds to the multiplication of two circles

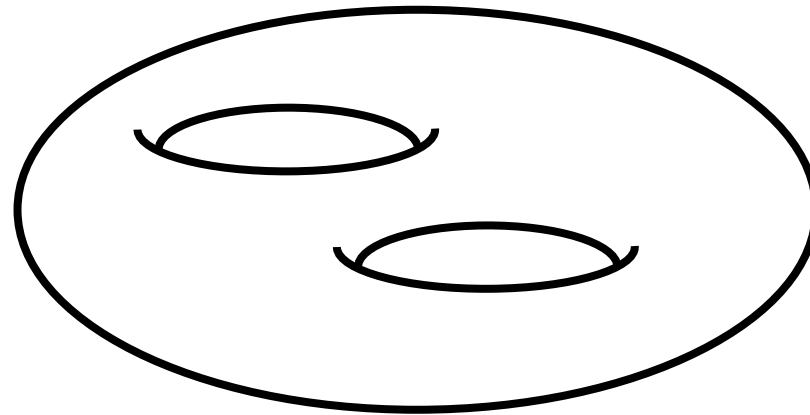


$$\mathbb{T} = \bigcup_{x \in S^1} S^1_x \approx S^1 \times S^1$$

$$\pi_1(\mathbb{T}) = \mathbb{Z} \times \mathbb{Z}$$

Topological invariant: (n_1, n_2)

Bonus question: Homotopy group of a sphere with many holes



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Projection of the geometrical space over the order parameter space

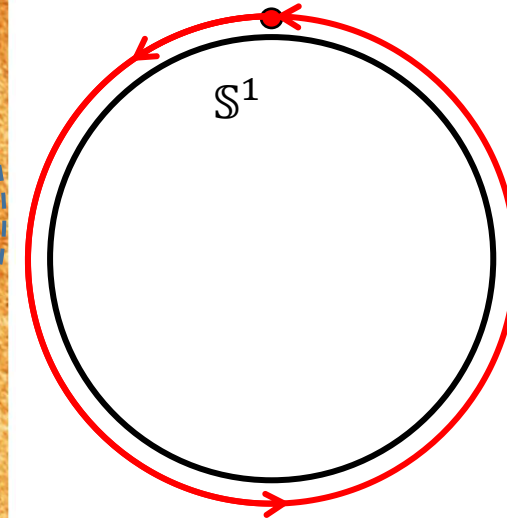
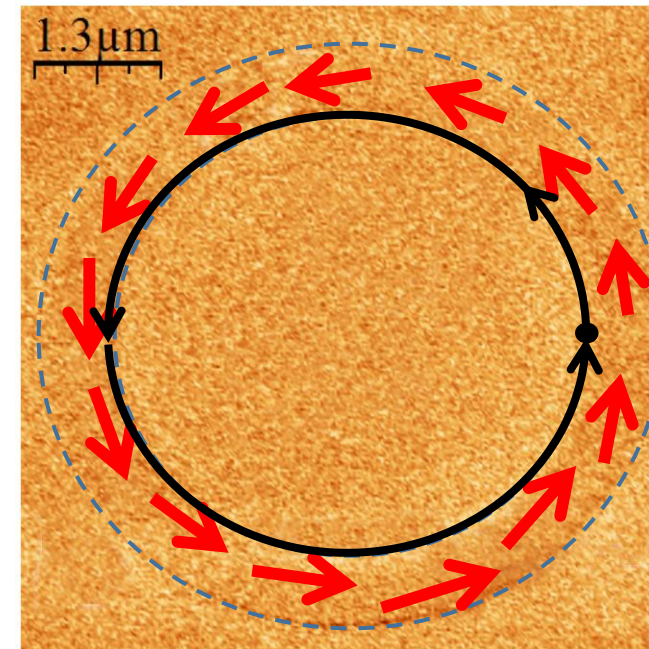
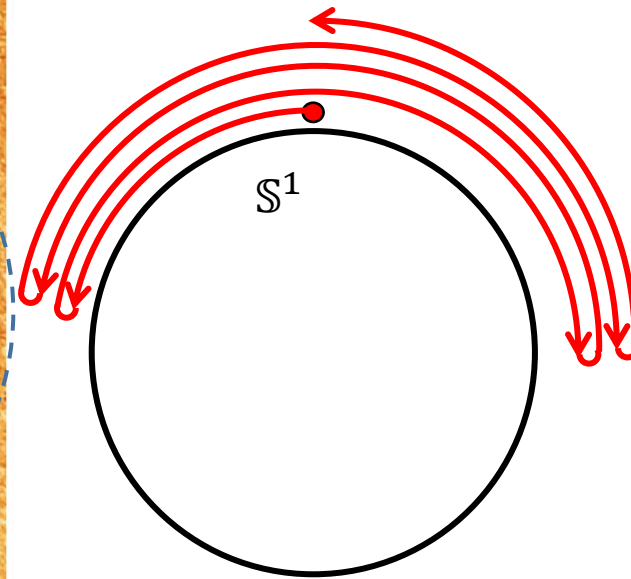
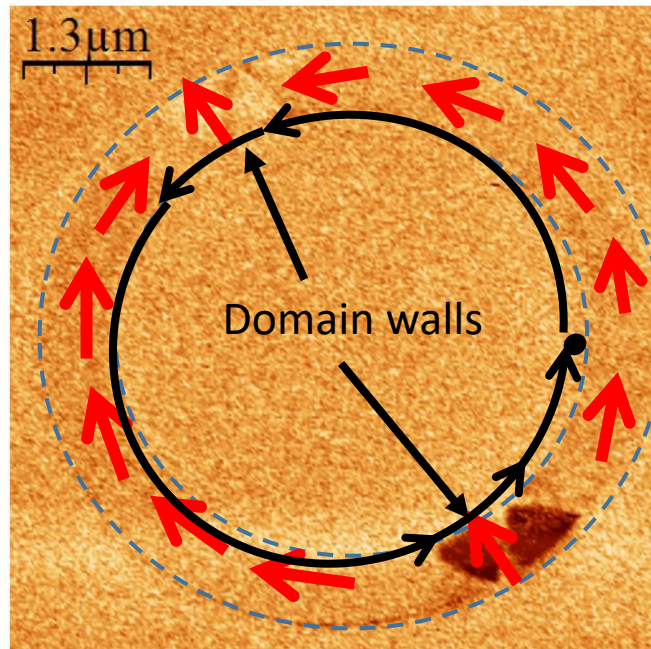
1. Easy situation: the sample has the same geometry as the order parameter space.

Ex: soft magnetic ring

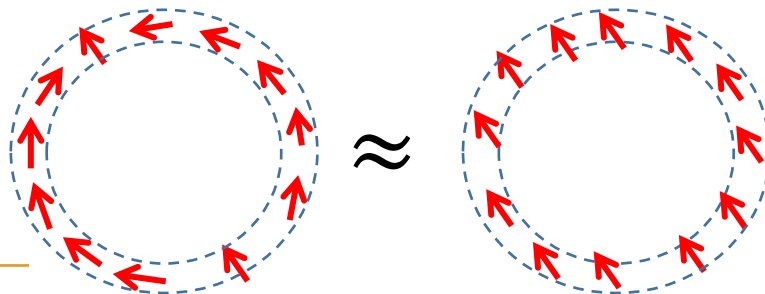
The dipolar interaction forces the magnetization along the ring ($\vec{m} \cdot \vec{n} = 0$).

Magnetization is S^1 -like

Magnetic force microscopy images of NiFe magnetic ring
J.Y. Chauleau (LPS, 2011)



$n = 0$

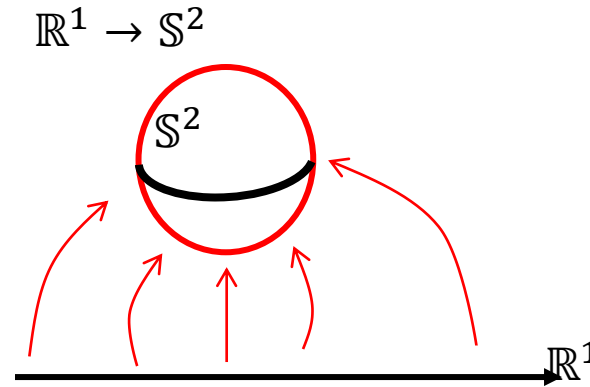
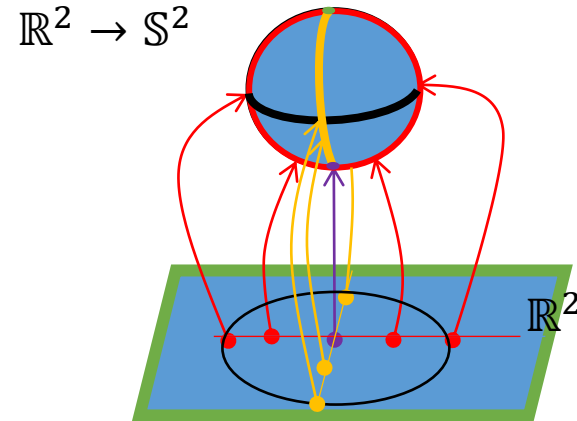
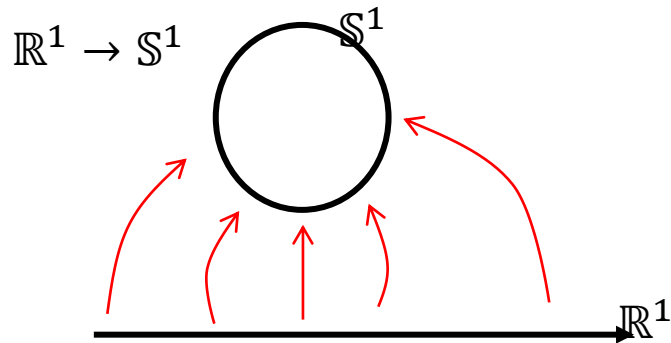


$n = 1$

Projection of the geometrical space over the order parameter space

2. The geometrical and order parameter spaces are different

Isomorphism is possible if boundary conditions are uniform



$\mathbb{R}^3 \rightarrow S^2$???

⇒ Follow the path and look to its trace on the order parameter space.

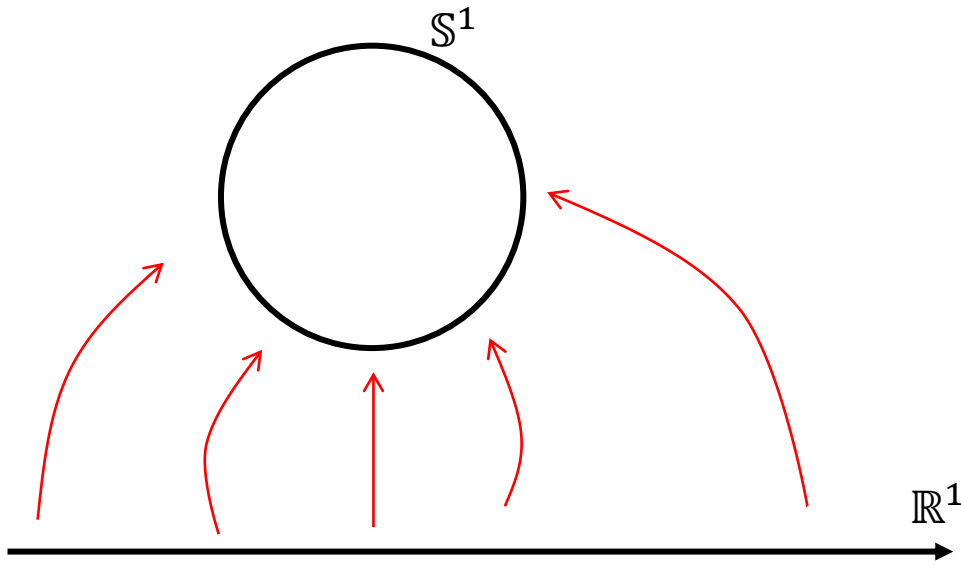
⇒ The trace is a contour of the same dimension as the geometrical space

Look for $\pi_d(S^p)$ homotopy group

- d : geometrical space dimension
- p : order parameter space dimension

$$\mathbb{R}^1 \rightarrow \mathbb{S}^1$$

We consider a 1D space \mathbb{R}^1 with some spins \mathbb{S}^1
 \Rightarrow We look to the trace of the order parameter on \mathbb{S}^1 when we move along \mathbb{R}^1 .



It corresponds to the developement performed to find the topological invariant of \mathbb{S}^1 .

The homotopy group is analogue to the homotopy group $\pi_1(\mathbb{S}^1) = \mathbb{Z}$

Example: 360° domain wall



The topological invariant is the (signed) number of 360° turns

If $\vec{S}(x) = (\cos \theta(x), \sin \theta(x))$

$$n = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \frac{d\theta}{dx}(x) dx = \frac{1}{2\pi} \int_{\theta(-\infty)}^{\theta(+\infty)} d\theta$$

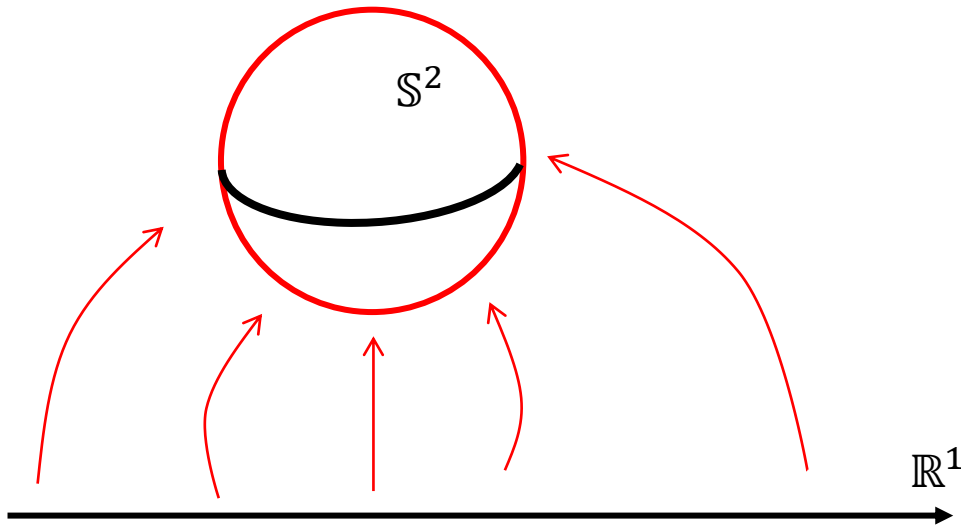
A 360° domain wall in \mathbb{S}^1 cannot collapse
 Collapse require new degree of freedom
 (change the norm of the spin or \mathbb{S}^2 spins)

$$\mathbb{R}^1 \rightarrow \mathbb{S}^2$$

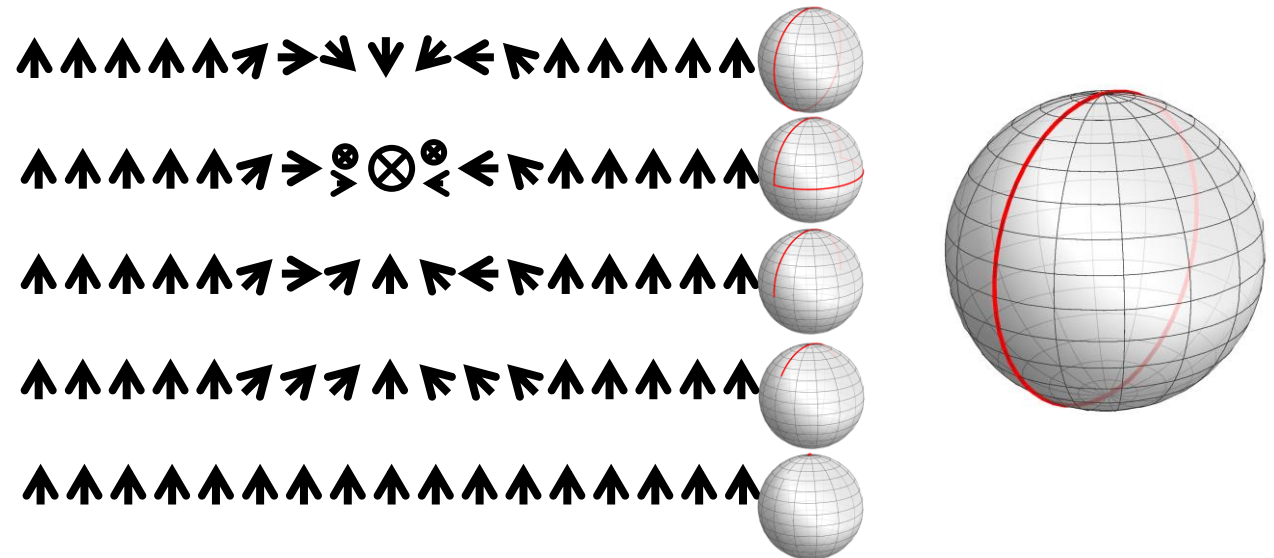
Same work using Heisenberg spins \mathbb{S}^2
 \Rightarrow We look to the trace of the order parameter on \mathbb{S}^2 when we move along \mathbb{R}^1 .

The order parameter space is partially covered
 The homotopy group is analogue to the homotopy group $\pi_1(\mathbb{S}^2) = 0$

Example: 360° domain wall



Collapse path to trivial state

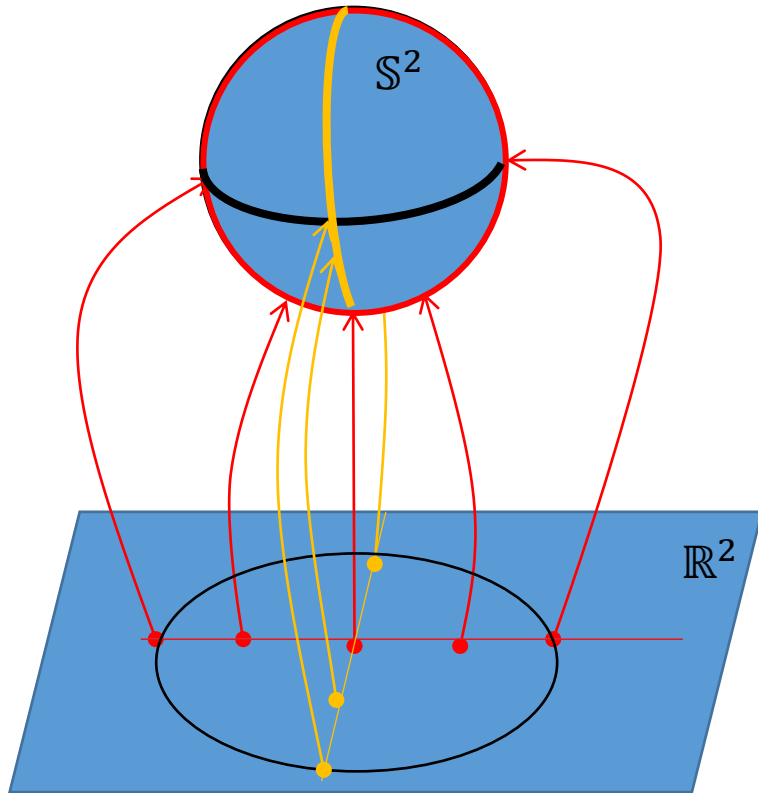


\Rightarrow Adding a new degree of freedom removes the topological protection

$$\mathbb{R}^2 \rightarrow \mathbb{S}^2$$

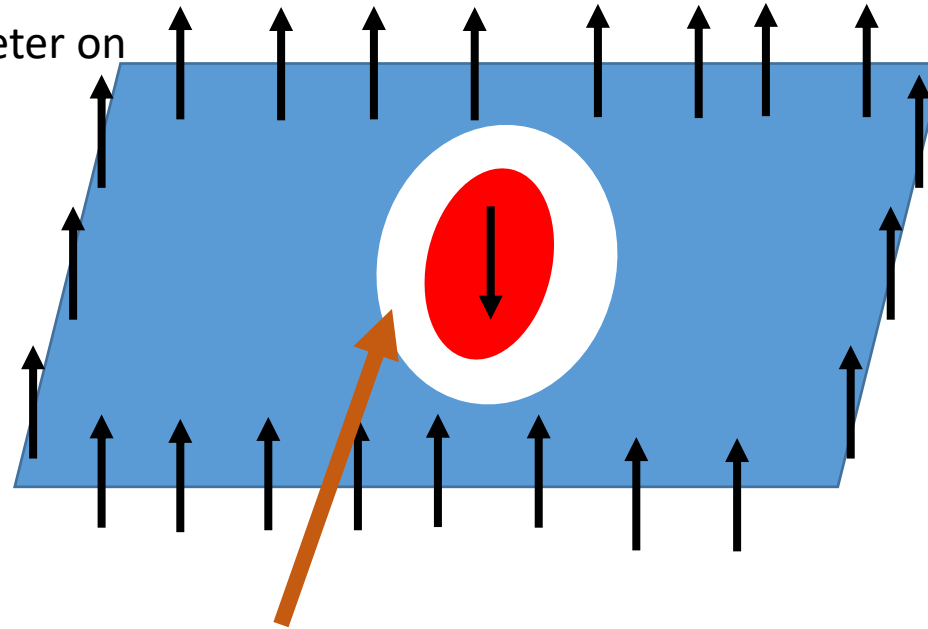
2D space using Heisenberg spins \mathbb{S}^2

\Rightarrow We look to the trace of the order parameter on \mathbb{S}^2 when we move along \mathbb{R}^2 .

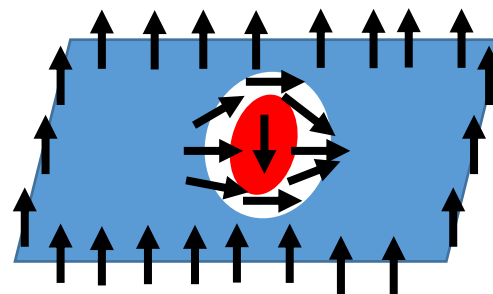
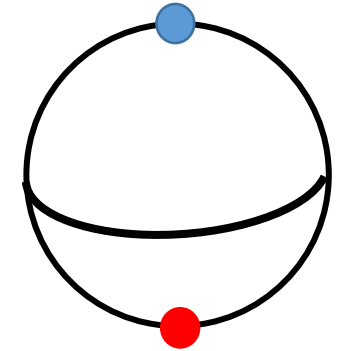


The homotopy group is analogue to the homotopy group $\pi_2(\mathbb{S}^2) = \mathbb{Z}$

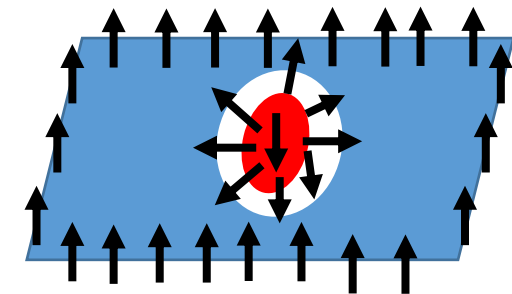
Example: Magnetic bubbles and skyrmions



The key is in the transition zone

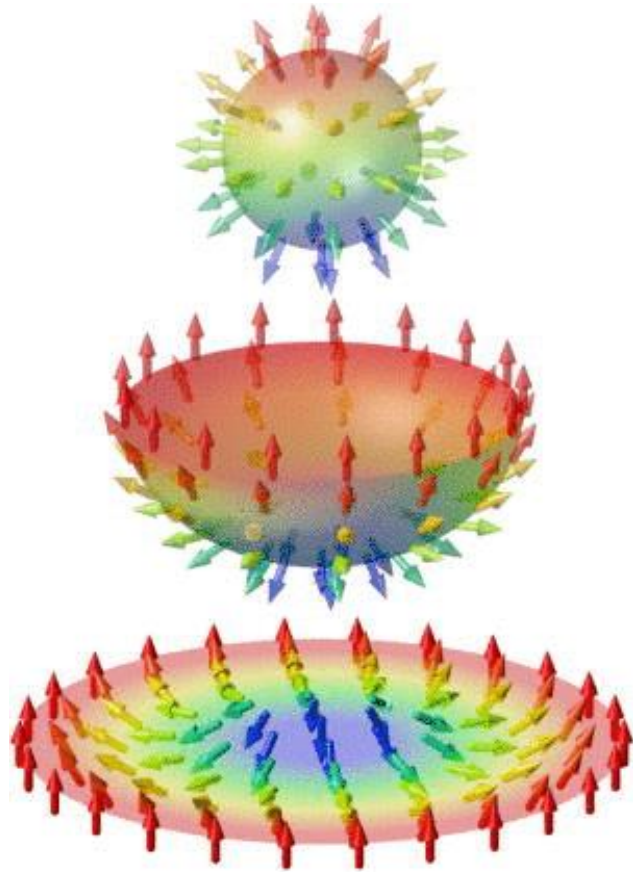


Non topological solution
 $n = 0$



Topological solution
 $n = 1$

$\mathbb{R}^2 \rightarrow \mathbb{S}^2$: Magnetic skyrmions



The topological invariant is the number of times the unit sphere is covered

$$n = \frac{1}{4\pi} \iint \vec{m} \cdot \left(\frac{\partial \vec{m}}{\partial x} \times \frac{\partial \vec{m}}{\partial y} \right) d^2r$$

For a centrosymmetric texture:

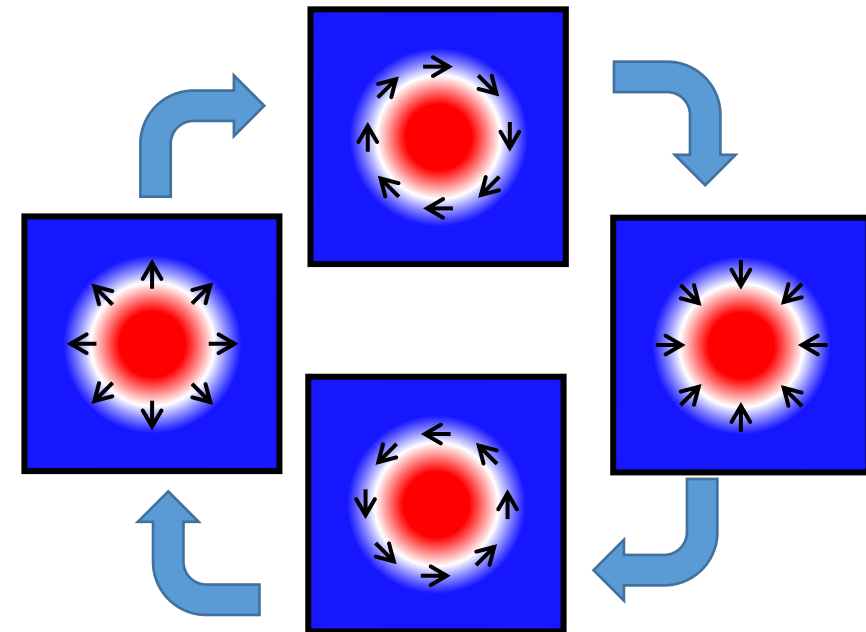
$$n = \frac{1}{4\pi} \int \sin \theta \frac{d\theta}{dr} dr \int \frac{d\phi}{d\varphi} d\varphi = pW$$

$$2p = m_z(r=0) - m_z(r=\infty)$$

$p \in \{-1, 0, 1\}$

Winding number W
 $W \in \mathbb{Z}$
 (topological invariant in $\pi_1(\mathbb{S}^1)$)

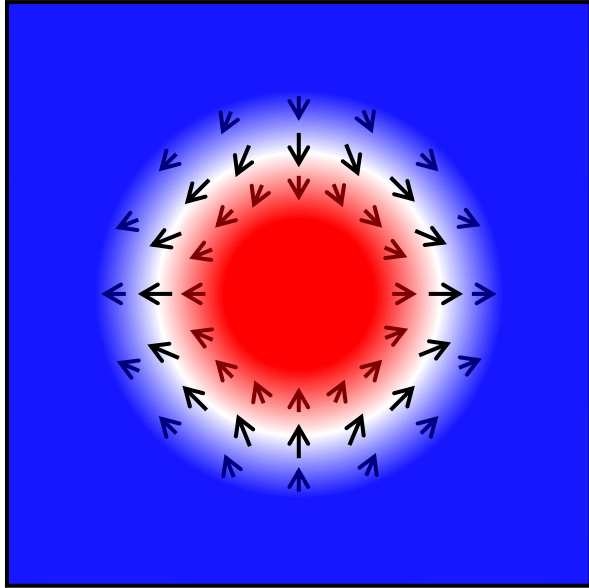
- Topology does not depend on the chirality
- Néel and Bloch skyrmions belong to the same homotopy class



$\mathbb{R}^2 \rightarrow \mathbb{S}^2$: Magnetic skyrmions and topological elements

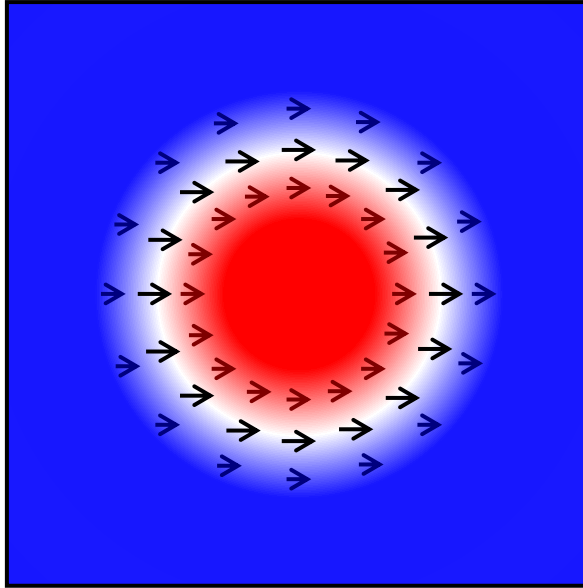
Antiskyrmion

$$p = 1; W = -1$$
$$n = -1$$



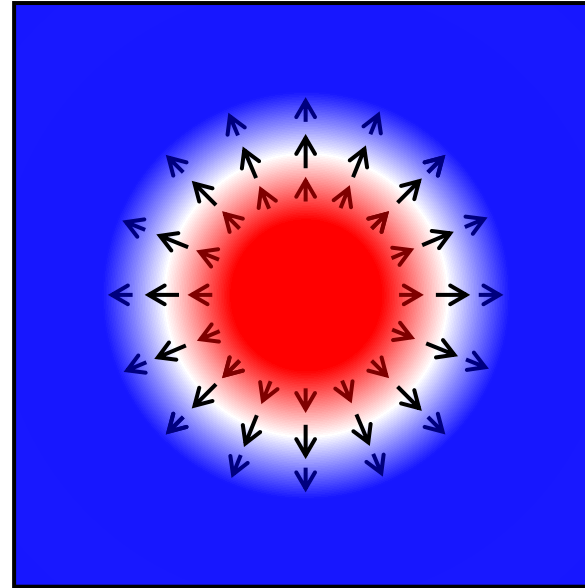
Trivial element:

$$p = 1; W = 0$$
$$n = 0$$



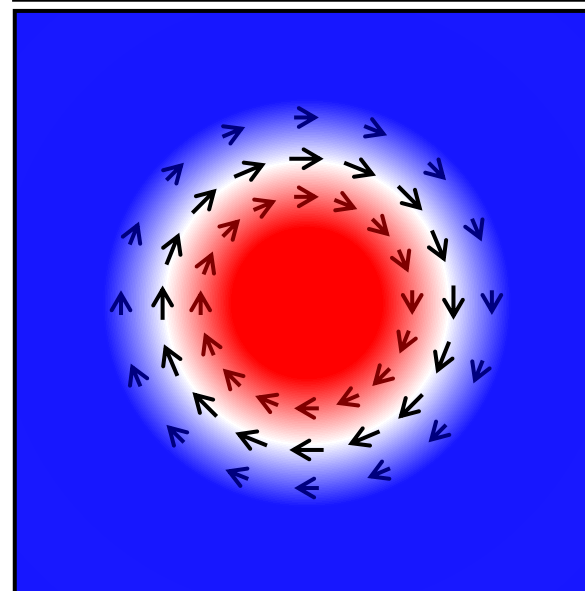
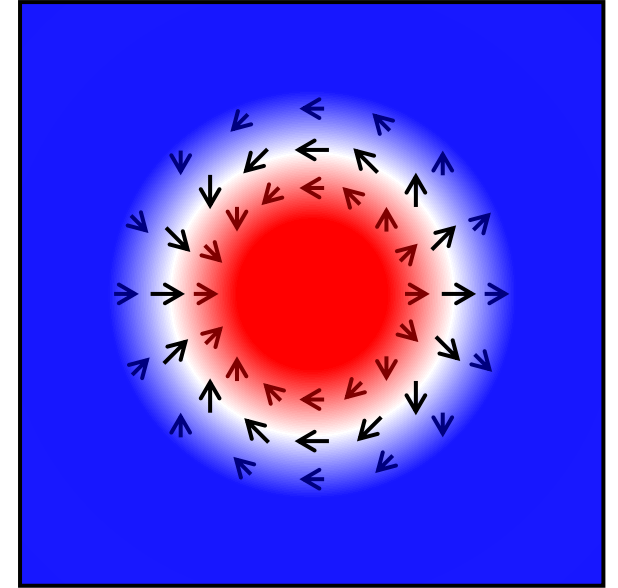
Skyrmion

$$p = 1; W = 1$$
$$n = 1$$



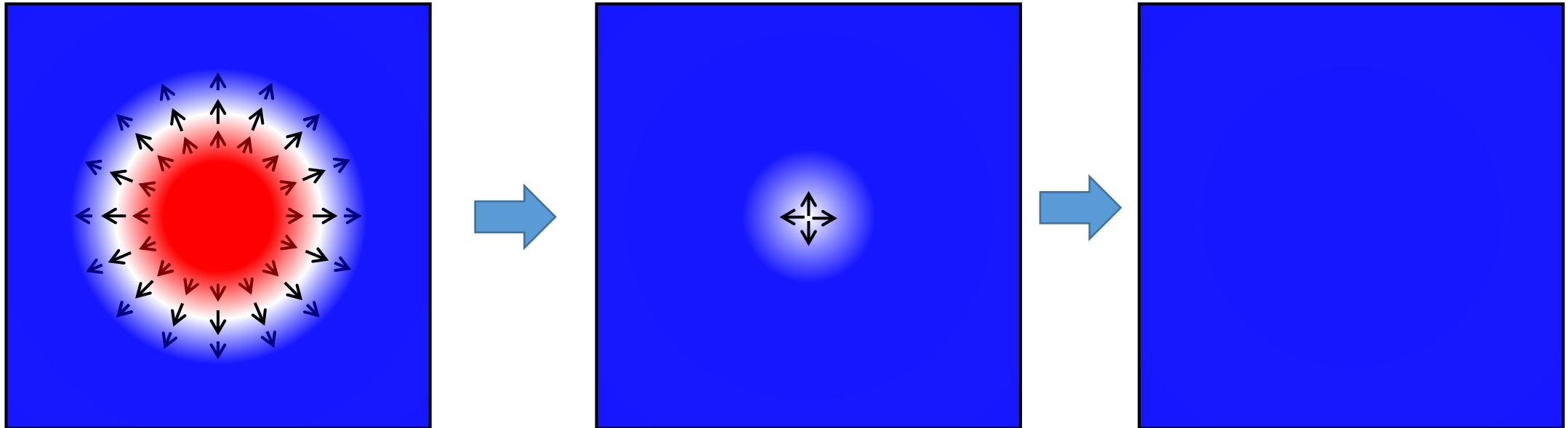
High order skyrmion:

$$p = 1; W = 2$$
$$n = 2$$



$\mathbb{R}^2 \rightarrow \mathbb{S}^2$: Magnetic skyrmions and topological elements

Collapse of a skyrmion



The collapse requires a topological defect

- Vortex for a 2D skyrmion

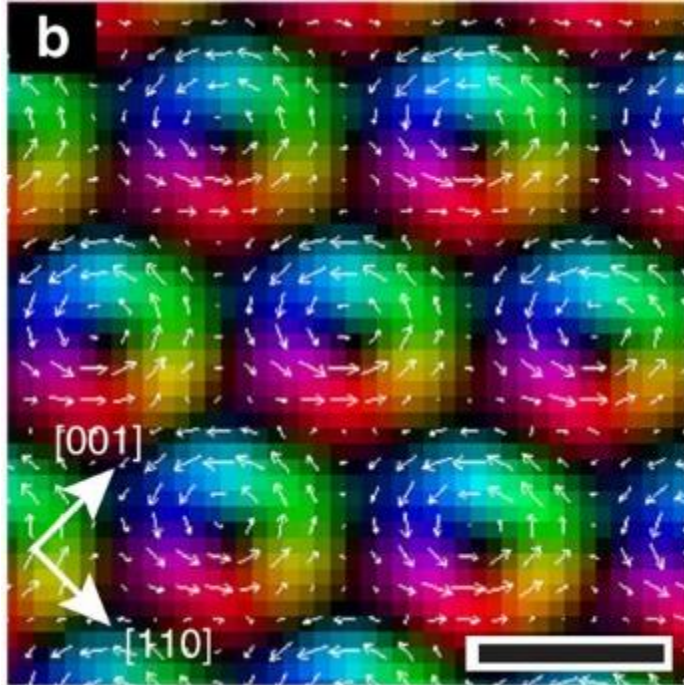
Topology changes requires to locally cancel the order parameter

⇒ Escape of 2D sphere \mathbb{S}^2 on the 2D ball \mathbb{B}^2 [$\pi_2(\mathbb{B}^2) = 0$]

$\mathbb{R}^2 \rightarrow \mathbb{S}^2$: Magnetic skyrmions

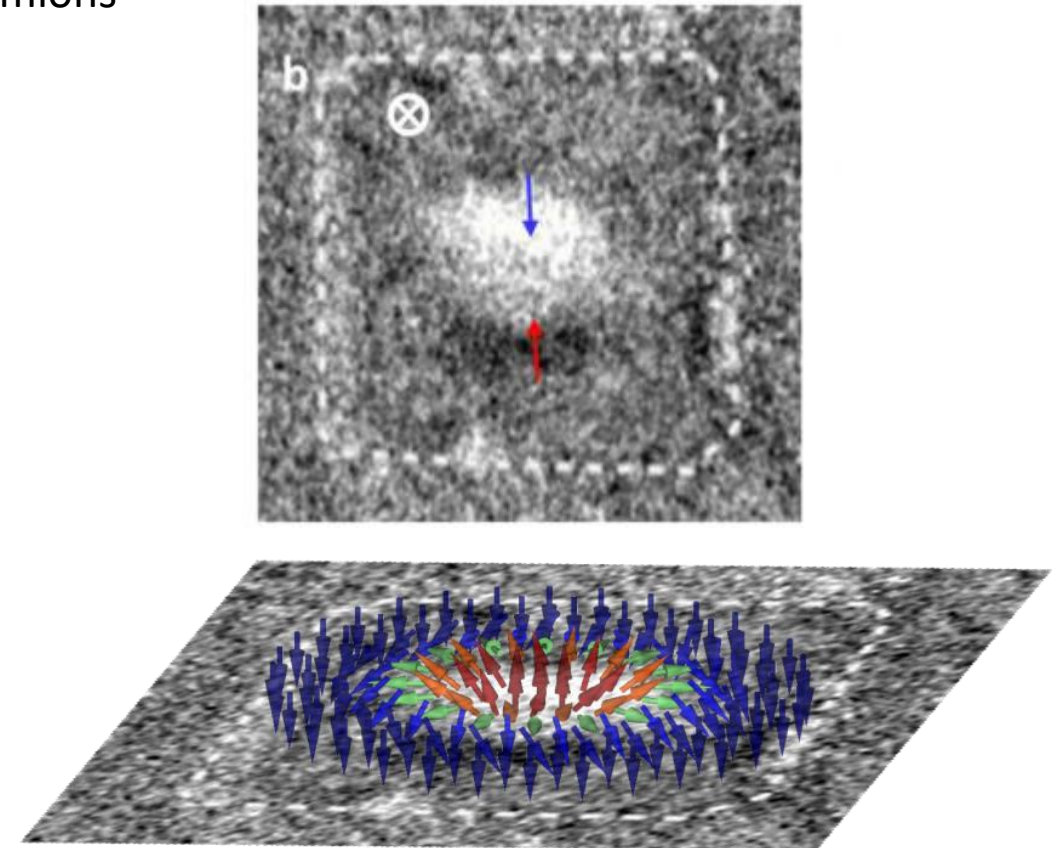
Skyrmions are obtained in situations with an absence of inversion symmetry
Dzyaloshinskii Moriya interaction fixes the chirality and therefore the topology

In Bulk crystals: Bloch skyrmions



[Yu et al. Nature (2010)]

In ultrathin films with structural absence of inversion symmetry:
Néel skyrmions



[Boule et al. Nature Nano 2016]

$\mathbb{R}^3 \rightarrow \mathbb{S}^2$: Hopfions

Space dimension is smaller than order parameter space:

\Rightarrow Isospins are lines in \mathbb{R}^3 .

Uniform condition at infinity

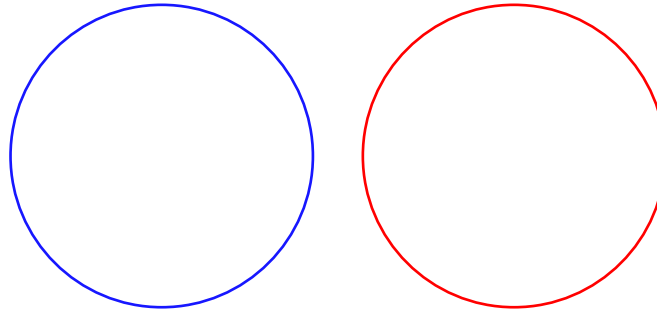
\Rightarrow Isospins are circles

$\Rightarrow \mathbb{R}^3 \approx \mathbb{S}^2 \times \mathbb{S}^1$

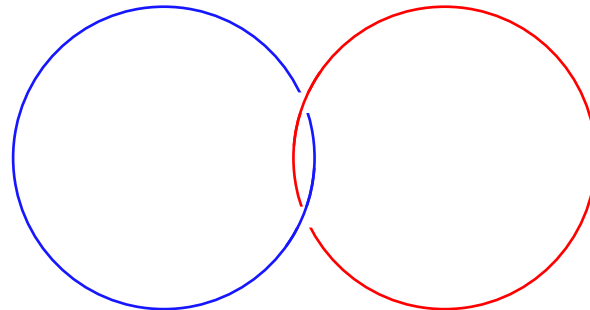
Homotopy group $\pi_3(\mathbb{S}^2)$?

$$\pi_3(\mathbb{S}^2) = \mathbb{Z}$$

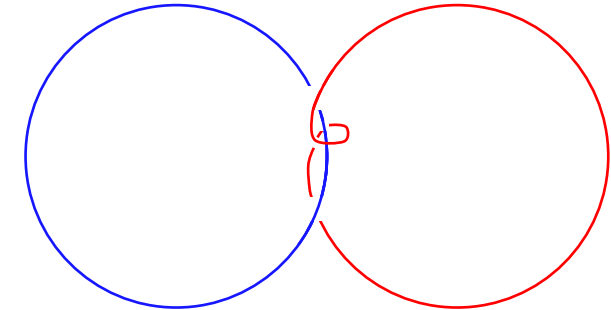
Trivial situation



Non trivial situations

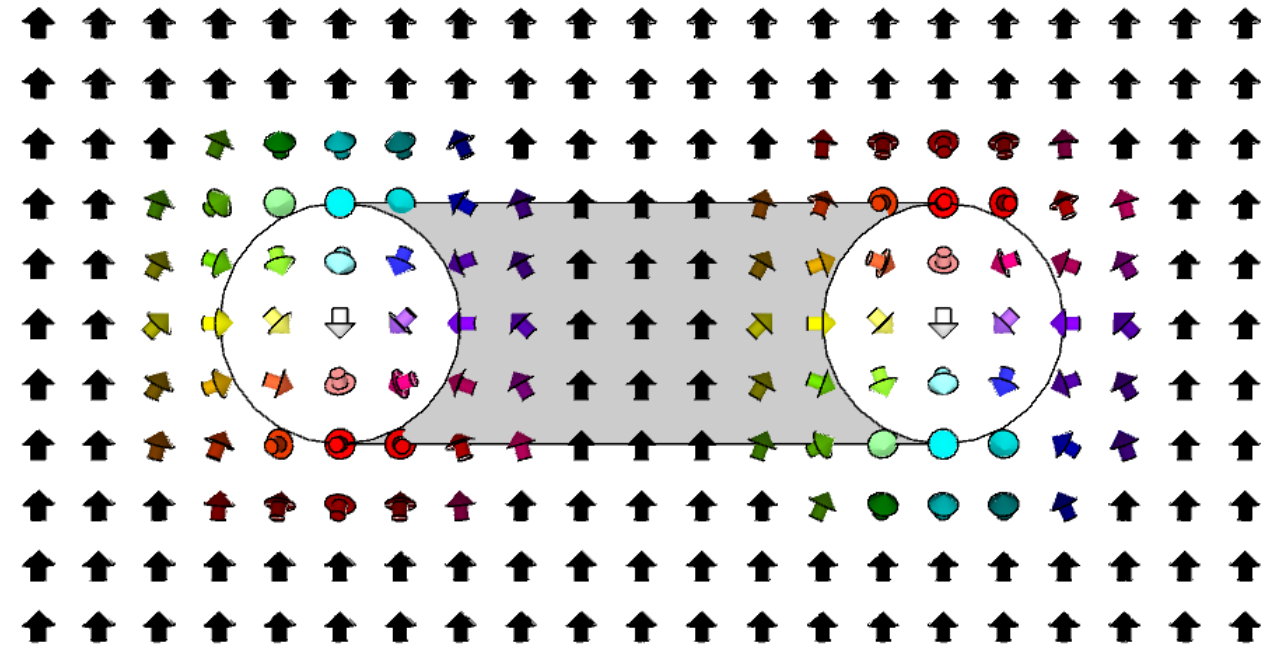
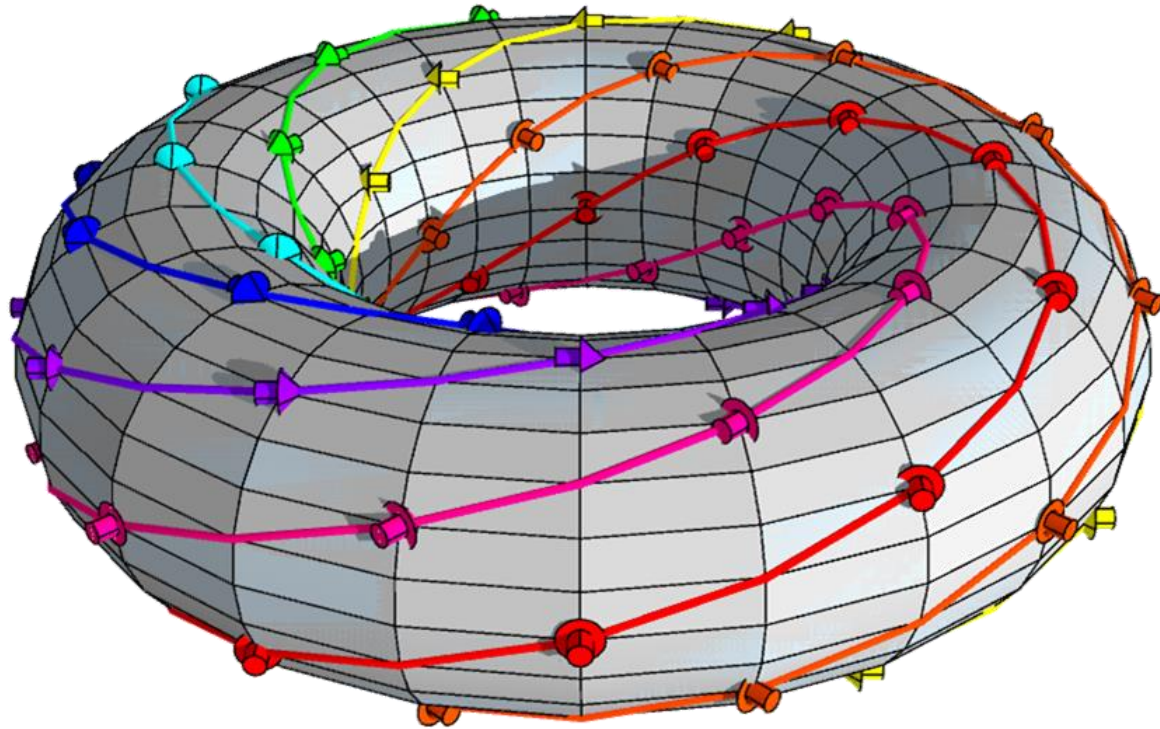


$n = 1$



$n = 2$

$\mathbb{R}^3 \rightarrow \mathbb{S}^2$: Hopfions



Projection of the space over the order parameter space

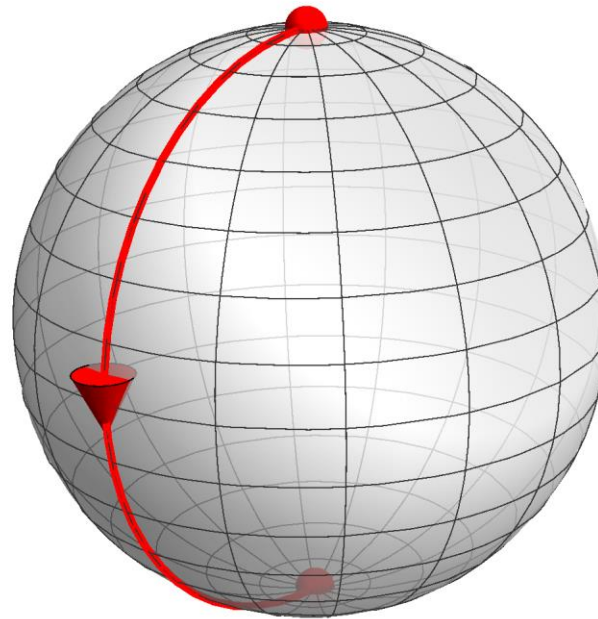
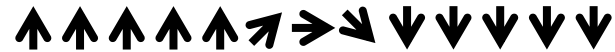
Case of non-uniform boundary conditions

The space **partially** wraps the order parameter space

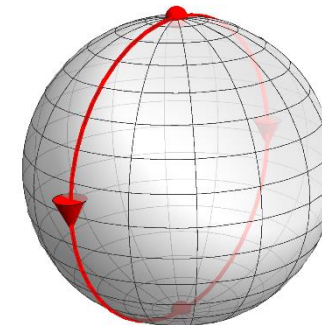
Boundary conditions are fixed by some energy => new topological solution

-1D: domain wall

Topological even with Heisenberg spins



For XY spins (e.g. strong DMI, in plane anisotropy),
topological difference between chiralities since
magnetization is bound to a single circle S^1



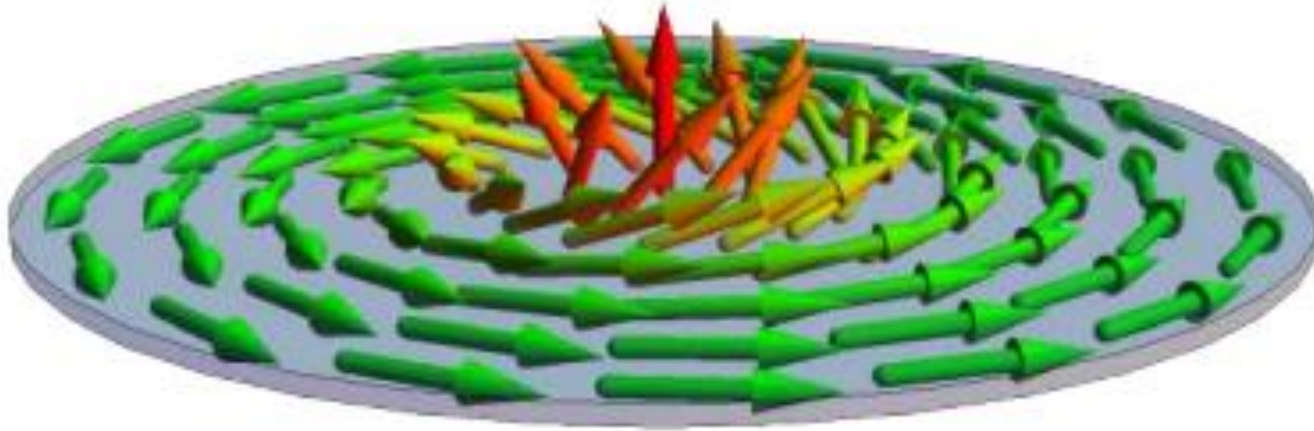
Projection of the space over the order parameter space

Case of non-uniform boundary conditions

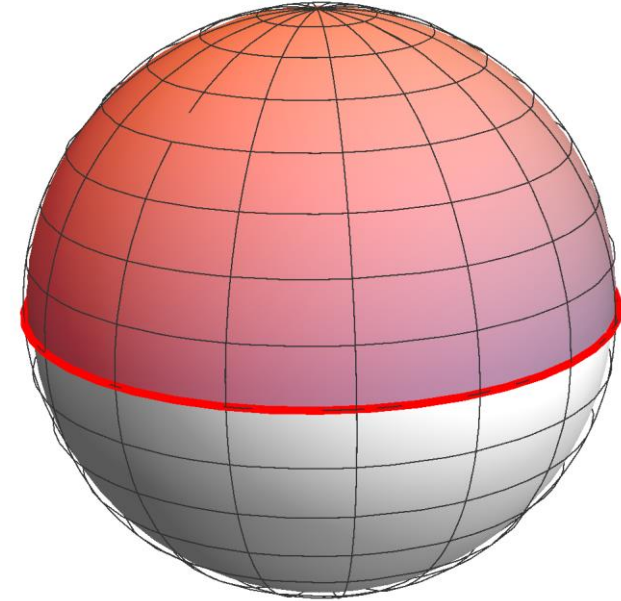
-2D space with Heisenberg spins: micromagnetic vortex (or meron)

-> Boundary condition: $\vec{m} \cdot \vec{n} = 0$ (magnetostatic charge minimization)

-> Dot center: $\vec{m} = \pm \vec{z}$ (exchange energy minimization)



(image from Benjamin Pigeau, Inst. Néel)



Topology depends on the vortex core orientation p and vorticity W

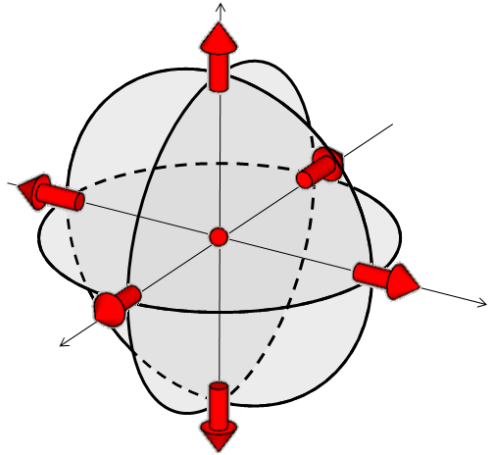
$$n = pW = \pm \frac{1}{2}$$

- 1. Introduction: homeomorphism**
- 2. How to catch a topological defect or texture**
- 3. Homotopy and homotopy group**
- 4. Geometrical space and order parameter space**
- 5. Topological defects and topologically stable configurations**

Concluding remark: topological defects and topologically stable structures

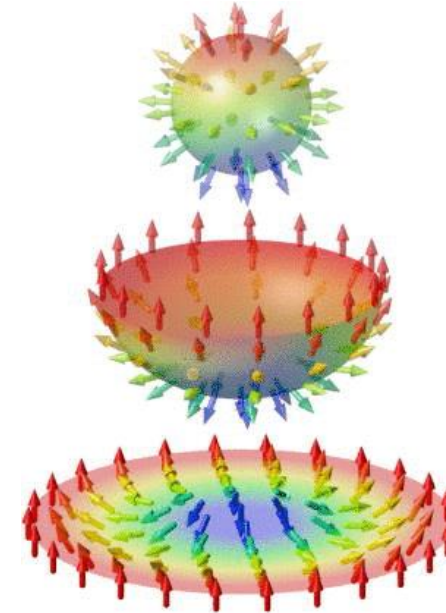
Ex: in $\pi_2(\mathbb{S}^2)$

Bloch point: topological defect



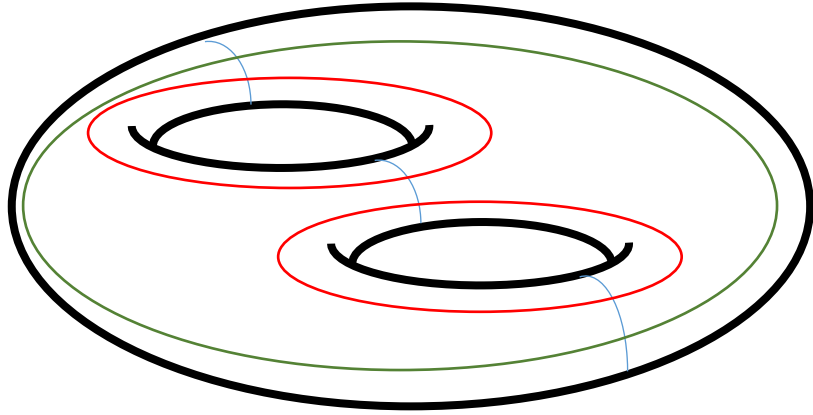
- Order parameter not defined at the defect
- Vanishing size
- Exchange energy diverges
- Stable by itself: is destroyed by a defect of opposite topological number

Skyrmion: topologically stable structure



- Order parameter always defined
- Finite size
- Finite exchange energy
- Can collapse through a topological defect

Homotopy group with n holes.



Genre 1 (tore): \mathbb{Z}^2

Genre 2: \mathbb{Z}^6

Genre 3: \mathbb{Z}^{13}

Genre 4: \mathbb{Z}^{29}

Genre 5: \mathbb{Z}^{61}